



## Time-dependent density functional theory study of the excited-state dihydrogen bond O–H···H–Si

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### ABSTRACT

Intermolecular dihydrogen bonding in the electronically excited states of a phenol–diethylmethylsilane (DEMS) complex was studied theoretically using the time-dependent density functional theory (TDDFT) method. Analysis of the frontier molecular orbitals revealed a locally excited  $S_1$  state for the dihydrogen-bonded phenol–DEMS complex in which only the phenol moiety is electronically excited. The calculated infrared spectrum of the phenol–DEMS complex is quite different from that of previously studied  $S_1$  state of a dihydrogen-bonded phenol–borane–trimethylamine complex. The O–H and Si–H stretching vibrational modes appear as intense, sharp peaks for the  $S_1$  state which are slightly red-shifted compared with those predicted for the ground state. Upon electronic excitation to the  $S_1$  state, the O–H and Si–H bonds involved in the dihydrogen bond O–H···H–Si lengthen slightly, while the C–O bond shortens. The calculated H···H distance is significantly shorter in the  $S_1$  state than in the ground state. Thus, the intermolecular dihydrogen bond of the phenol–DEMS complex is stronger in the electronically excited state than in the ground state.

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### 1. Introduction

Over the last decades, the hydrogen bond has gained wide attention because of its importance in physical, chemical and biological processes [1–51]. The conventional hydrogen bond is an analogous interaction  $X-H\cdots A$  formed by a strongly polar group  $X^{\delta-}-H^{\delta+}$  on one side, and atom  $A^{\delta-}$  with a lone pair of electrons on the other side [1–22]. As the study of hydrogen bonds has progressed, some unconventional hydrogen bonds have been proposed using diverse experimental techniques and theoretical calculations. Another type of hydrogen bond, formed between two oppositely charged hydrogen atoms, is called the dihydrogen bond [23–51]. It can be represented as  $X-H^{\delta-}\cdots\delta^+H-Y$ , where X and Y are electropositive and electronegative atoms respectively [23–29]. The electrostatic attraction between oppositely polarized hydrogen atoms promotes the formation of the dihydrogen bond and maintains its stability [25–35]. With comparative strength and directionality to conventional hydrogen bonding, dihydrogen bonding can influence the structure, reactivity and selectivity of compounds in gas and condensed phases. Thus, it has potential application in catalysis, crystal engineering and materials chemistry [23–49].

Dihydrogen bonding is now a well-established interaction and found in various complexes. Mikami and co-workers have reported the formation of several dihydrogen-bonded complexes between borane amines and molecules containing acidic hydrogen atoms in the gas phase [36–45]. Recently, Zhao and Han [49] reported the theoretical study on the structure and dynamics of a novel dihydrogen bond in the electronically excited state of a phenol–borane–trimethylamine (BTMA) complex for the first time. In their landmark work, their calculated results reveal that the dihydrogen-bonded phenol–BTMA complex possesses a locally excited (LE)  $S_1$  state that is centered on the phenol moiety. In addition, there is no O–H or B–H stretching vibrational mode in the calculated infrared (IR) spectrum of the  $S_1$  state of the dihydrogen-bonded phenol–BTMA complex. This was attributed to the strengthening of the dihydrogen bond O–H···H–B in the electronically excited state.

Ishikawa et al. [47] reported the first experimental observation of a dihydrogen bond involving a Si–H group in a phenol–diethylmethylsilane (DEMS) complex. This complex should be a good precursor for  $H_2$  elimination reactions. IR–UV double-resonance spectroscopy was performed to identify the dihydrogen bond containing the Si–H group. When DEMS was added to gaseous phenol, many small bands appeared in the vicinity of the origin band attributed to the phenol monomer in the fluorescence excitation spectrum. Three isomers of the complex were identified by IR–UV population labeling spectroscopy and laser-induced fluorescence spectroscopy. The electronic ori-

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gin band of the phenol–DEMS complex shifts by  $-78$  to  $+4$   $\text{cm}^{-1}$  compared with that of the phenol monomer. The IR absorption spectra of the phenol–DEMS isomers show that the O–H stretching vibration frequencies shift to lower frequency by  $21$ – $29$   $\text{cm}^{-1}$ . The magnitudes of these shifts are much smaller than those calculated for the phenol–borane–dimethylamine ( $-174$   $\text{cm}^{-1}$ ) [38] and phenol–BTMA ( $-143$   $\text{cm}^{-1}$ ) complexes [39]. To obtain structural information about the dihydrogen-bonded phenol–DEMS complex, Ishikawa et al. performed quantum-chemical calculations together with experimental measurements. The calculation showed there were five stable isomers in the ground state. For three isomers, the distance between the two hydrogen atoms involved in the dihydrogen bonds ( $R_{\text{H}\cdots\text{H}}$ ) is  $2.23$  Å, and the O–H vibrational frequency ( $\nu_{\text{OH}}$ ) is  $3641$   $\text{cm}^{-1}$ . The  $R_{\text{H}\cdots\text{H}}$  and  $\nu_{\text{OH}}$  of the other two isomers are  $2.27$  Å and  $3640$   $\text{cm}^{-1}$ , and  $2.27$  Å and  $3641$   $\text{cm}^{-1}$ , respectively. Further spectroscopic investigations and theoretical calculations are still required for various dihydrogen-bonded complexes to fully understand the properties of the dihydrogen bond, such as geometric conformation and its potential utility. Tremendous effort has been spent investigating the ground state properties of dihydrogen bonds in dihydrogen-bonded complexes. In contrast, research focused on the structure and dynamics of dihydrogen bonds in electronically excited states is scarce. In this study, theoretical calculations have been performed to investigate the properties of the dihydrogen bond in an electronically excited state of a dihydrogen-bonded phenol–DEMS complex. The five isomers proposed by Ishikawa et al. for this complex have similar geometries. The geometry of isomer 1 reported by Ishikawa et al. was used in the present work as it possesses the lowest relative energy and has a  $R_{\text{H}\cdots\text{H}}$  of  $2.23$  Å. Time-dependent density functional theory (TDDFT) was used for excited-state geometry optimization and electronic transition calculations for the dihydrogen-bonded phenol–DEMS complex and the isolated monomers. TDDFT has been proven a valid tool for calculating the IR spectra of the electronically excited state of compounds [9–22]. Here, TDDFT is used to predict the excited-state IR spectra of both the O–H and Si–H stretching vibrational regions of the dihydrogen-bonded phenol–DEMS complex.

## 2. Computational methods

Density functional theory (DFT) and time-dependent density functional theory (TDDFT) were employed for our ground state and excited state computations, respectively [49–52]. The resolution-of-the-identity (RI) approximation was adopted in the purpose of improving efficiency without sacrificing accuracy of the results [53–55]. All geometry optimizations were carried out using BP-86 functional with triple zeta valence plus polarization (TZVP) basis set as well as corresponding auxiliary basis sets for resolution-of-the-identity (RI) approximation [56]. Harmonic vibrational frequency analysis was performed at the same theoretical level [54–57]. All the computations were performed using the TURBOMOLE program suite.

## 3. Results and discussion

The optimized geometry of the dihydrogen-bonded phenol–DEMS complex in the ground state is shown in Fig. 1. In the ground state, the dihydrogen bond O–H<sub>1</sub>⋯H<sub>2</sub>–Si is formed between phenol and DEMS. Table 1 lists the calculated bond lengths, angles and dihedral angle of isolated phenol and DEMS and the dihydrogen-bonded phenol–DEMS complex in the ground state. The calculated distance between the two hydrogen atoms H<sub>1</sub>⋯H<sub>2</sub> is  $1.841$  Å, which is slightly longer than that of the dihydrogen-bonded phenol–BTMA complex ( $1.737$  Å) [49]. The bond angles O–H<sub>1</sub>⋯H<sub>2</sub> and H<sub>1</sub>⋯H<sub>2</sub>–Si for the optimized con-

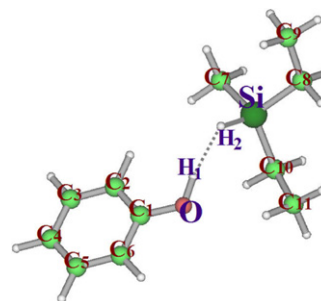


Fig. 1. Optimized geometry of the dihydrogen-bonded phenol–DEMS complex in the ground state. The dotted line denotes the intermolecular dihydrogen bond.

formation of the phenol–DEMS complex are  $170.8^\circ$  and  $138.3^\circ$ , respectively. The dihedral angle O–H<sub>1</sub>⋯H<sub>2</sub>–Si is  $46.65^\circ$ . Furthermore, the calculated bond length for the O–H<sub>1</sub> group is  $0.976$  Å in the dihydrogen-bonded phenol–DEMS complex, which is slightly longer than that of the O–H<sub>1</sub> group in isolated phenol ( $0.973$  Å). At the same time, the bond length for C<sub>1</sub>–O in the phenol–DEMS complex ( $1.360$  Å) is slightly shorter than that in free phenol ( $1.376$  Å). In the DEMS moiety, the Si–H<sub>2</sub> bond lengths and the Si–C<sub>7,8,10</sub> bonds shorten upon formation of the intermolecular dihydrogen bond. These results are similar to those observed for the dihydrogen-bonded phenol–BTMA complex.

The electronic transition energies and corresponding oscillator strengths for the low-lying electronically excited states of the dihydrogen-bonded phenol–DEMS complex and the isolated monomers were calculated using TDDFT. The results are listed in Table 2. All of the electronic transition energies of the dihydrogen-bonded complex are just a little higher than those of isolated phenol, but significantly lower than those of free DEMS. The phenol–DEMS complex exhibits phenol-based elec-

Table 1

Calculated bond lengths (Å), angles ( $^\circ$ ) and dihedral angle ( $^\circ$ ) of the isolated monomers and the dihydrogen-bonded phenol–DEMS complex in the ground state.

	Phenol <sup>a</sup>	DEMS	Phenol–DEMS
C <sub>1</sub> –O	1.376	–	1.360
O–H <sub>1</sub>	0.973	–	0.976
H <sub>1</sub> ⋯H <sub>2</sub>	–	–	1.841
H <sub>2</sub> –Si	–	1.506	1.514
O–H <sub>1</sub> ⋯H <sub>2</sub>	–	–	170.8
H <sub>1</sub> ⋯H <sub>2</sub> –Si	–	–	138.3
O–H <sub>1</sub> ⋯H <sub>2</sub> –Si	–	–	46.65
Si–C <sub>7</sub>	–	1.896	1.893
Si–C <sub>8</sub>	–	1.904	1.900
Si–C <sub>10</sub>	–	1.903	1.899

<sup>a</sup> Data taken from Ref. [49].

Table 2

Calculated electronic transition energies (eV) and corresponding oscillator strengths (in parentheses) for the low-lying electronically excited states of the dihydrogen-bonded phenol–DEMS complex and the isolated monomers. The contributions of the orbital transitions for the S<sub>1</sub> state are also listed.

	Phenol–DEMS	Phenol <sup>a</sup>	DEMS
S <sub>1</sub>	5.114(0.033)	4.789(0.031)	8.010(0.008)
	H → L 80.2%	H → L 82.5%	H → L 91.3%
	H-1 → L+1 19.3%	H-1 → L+1 16.7%	
S <sub>2</sub>	5.967(0.058)	5.684(0.049)	8.134
S <sub>3</sub>	6.208(0.000)	5.782(0.000)	8.336
S <sub>4</sub>	6.509(0.003)	6.467(0.001)	8.486
S <sub>5</sub>	6.812(0.208)	6.508(0.215)	8.514
S <sub>6</sub>	6.944(0.519)	6.575(0.002)	8.746

H: highest occupied molecular orbital (HOMO), L: lowest unoccupied molecular orbital (LUMO).

<sup>a</sup> Data taken from Ref. [49].

tronic transitions that are blue-shifted compared with that of free phenol, i.e., the electronic transitions of phenol increase in energy upon formation of the intermolecular dihydrogen bond. The electronic transition energy for the  $S_1$  state of the dihydrogen-bonded phenol–DEMS complex (5.114 eV) is close to that of the phenol monomer (4.789 eV), which is much lower than the transition energy for  $S_1$  state of isolated DEMS (8.010 eV). This infers that the electronically excited state of the dihydrogen-bonded phenol–DEMS complex is localized on the phenol molecule, while the DEMS molecule remains in the ground state. The orbital transition contributions for the  $S_1$  state of the phenol–DEMS complex are also presented in Table 2. Two main orbital transitions contribute to the  $S_1$  state: one is from the highest occupied molecular orbital (HOMO) to the lowest unoccupied molecular orbital (LUMO) (80.2%), the other is from the HOMO-1 to the LUMO+1 (19.3%). Thus, it can be concluded that four main orbitals are involved in the  $S_1$  state of the dihydrogen-bonded phenol–DEMS complex. This finding varies from the  $S_1$  state of the dihydrogen-bonded phenol–BTMA complex which involves three molecular orbitals: the HOMO, LUMO and LUMO+1.

Molecular orbital (MO) analysis can directly provide insight into the nature of the excited states [58–64]. The frontier MOs of the dihydrogen-bonded phenol–DEMS complex are depicted in Fig. 2. Here, only the HOMO-1, HOMO, LUMO and LUMO+1 orbitals are shown because the  $S_1$  state of the complex is associated with these four orbitals. The electron densities of all four orbitals are localized on the phenol moiety. Hence, the  $S_1$  state of the dihydrogen-bonded phenol–DEMS complex is a LE state. This result is consistent with our earlier speculation that only the phenol moiety is excited in the  $S_1$  state of the dihydrogen-bonded phenol–DEMS complex. From Fig. 2, one can visualize electronic excitation between the  $\pi$ – $\pi^*$  orbitals of the benzene ring and the lone electron pair on the hydroxyl group. Thus, both the dihydrogen bond  $O-H_1 \cdots H_2-Si$  and the phenol moiety in the dihydrogen-bonded phenol–DEMS complex should be strongly influenced by charge redistribution in the  $S_1$  state, which is similar to that of the dihydrogen-bonded phenol–BTMA complex.

The excited-state geometry optimizations for the dihydrogen-bonded phenol–DEMS complex and the isolated monomers were performed using TDDFT. Based on the optimized structures, the IR spectra for both the ground state and the  $S_1$  state were calculated. Changes in the intermolecular dihydrogen bonding  $O-H_1 \cdots H_2-Si$  between the different states can be monitored by shifts in the O–H and Si–H stretching vibrational modes in the IR spectra. The calculated IR spectra for the dihydrogen-bonded phenol–DEMS complex and isolated monomers in different electronic states are shown in Fig. 3. All of the O–H and Si–H stretching vibration frequencies are labeled in Fig. 3. The O–H stretching vibration frequency of isolated

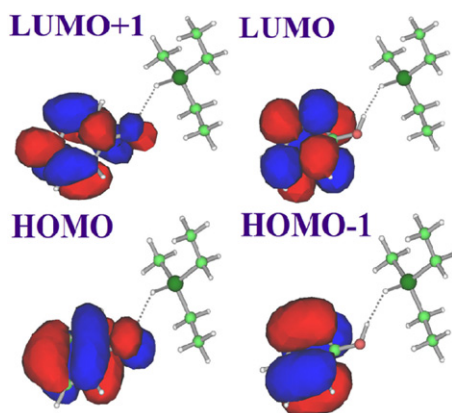


Fig. 2. Frontier MOs of the dihydrogen-bonded phenol–DEMS complex.

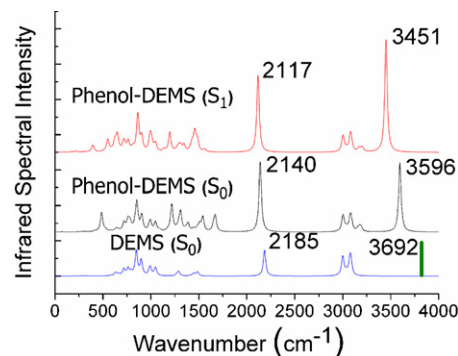


Fig. 3. Calculated IR spectra of the dihydrogen-bonded phenol–DEMS complex and DEMS in different electronic states. The thick green line denotes the stretching vibrational frequency of the hydroxyl group in isolated phenol [49]. (For interpretation of the references to color in this figure legend, the reader is referred to the web version of the article.)

phenol in the ground state ( $3692\text{ cm}^{-1}$ ) was calculated by Zhao and Han [49]. The calculated stretching vibration frequency for the O–H bond in the dihydrogen-bonded phenol–DEMS complex shifts to  $3596\text{ cm}^{-1}$  because of the formation of the dihydrogen bond  $O-H_1 \cdots H_2-Si$ . The magnitude of the shift ( $96\text{ cm}^{-1}$ ) is much smaller than that calculated for the dihydrogen-bonded phenol–BTMA complex ( $202\text{ cm}^{-1}$ ) [49]. This indicates that the dihydrogen bond  $O-H \cdots H-Si$  is not as strong as the dihydrogen bond  $O-H \cdots H-B$  in the ground state. The stretching vibration frequency at  $2140\text{ cm}^{-1}$  can be assigned to the Si–H vibrational mode of the complex, which is shifted to a lower frequency by  $45\text{ cm}^{-1}$  relative to free DEMS. Furthermore, in the  $S_1$  state, the O–H stretching vibrational mode of the dihydrogen-bonded phenol–DEMS complex shows a red shift of  $145\text{ cm}^{-1}$  compared with the ground state, indicating that the  $O-H_1$  bond lengthens upon electronic excitation. The Si–H group of the dihydrogen-bonded phenol–DEMS complex exhibits a relatively small red shift from the ground state ( $2140\text{ cm}^{-1}$ ) to the  $S_1$  state ( $2117\text{ cm}^{-1}$ ). Therefore, the length of the Si– $H_2$  bond may not increase as much as the  $O-H_1$  bond upon electronic excitation. The changes in the  $O-H_1$  and Si– $H_2$  bonds suggest that the dihydrogen bond  $O-H_1 \cdots H_2-Si$  should be stronger in the  $S_1$  state than the ground state. The IR signals of the O–H and Si–H groups involved in dihydrogen bonding are intense and sharp in the  $S_1$  state. This is quite different behavior from that of the  $S_1$  state of the dihydrogen-bonded phenol–BTMA complex. The stretching vibrational bands of the O–H and B–H bonds of the dihydrogen-bonded phenol–BTMA complex both disappeared in the  $S_1$  state, followed by a strong vibrational absorption appearing in the region of a C=O stretching vibration [49]. The differences between the excited state behavior of the two complexes suggest that the  $O-H_1$  and Si– $H_2$  bonds of the dihydrogen-bonded phenol–DEMS complex may be lengthened less than the O–H and B–H bonds of the dihydrogen-bonded phenol–BTMA complex upon electronic excitation.

Table 3 lists all the bond lengths, angles and the dihedral angles of the dihydrogen bond for different electronic states of the phenol–DEMS complex. The  $C_1-O$  bond shortens by  $0.016\text{ \AA}$  in the ground-state phenol–DEMS complex due to formation of the intermolecular dihydrogen bond. Upon electronic excitation to the  $S_1$  state in the dihydrogen-bonded phenol–DEMS complex, the  $C_1-O$  and  $O-H_1$  bonds change by  $-0.014$  and  $+0.009\text{ \AA}$ , respectively. These changes are much smaller than those found for the dihydrogen-bonded phenol–BTMA complex. These findings are consistent with the IR spectra of the dihydrogen-bonded phenol–DEMS complex, because the  $S_1$  state exhibits an O–H stretching band that does not show C=O bond character. For the DEMS moiety of the complex in ground state, the Si– $H_2$  bond lengthens slightly by  $0.008\text{ \AA}$  after formation of the intermolec-

**Table 3**

Calculated bond lengths (Å), angles (°) and dihedral angles (°) of the dihydrogen-bonded phenol–DEMS complex in the ground state (GS) and the electronically excited state (ES).

	Phenol–DEMS	
	GS	ES
C <sub>1</sub> –O	1.360	1.346
O–H <sub>1</sub>	0.976	0.985
H <sub>1</sub> ···H <sub>2</sub>	1.841	1.712
H <sub>2</sub> –Si	1.514	1.518
O–H <sub>1</sub> ···H <sub>2</sub>	170.8	171.5
H <sub>1</sub> ···H <sub>2</sub> –Si	138.3	132.0
O–H <sub>1</sub> ···H <sub>2</sub> –Si	46.65	76.32
Si–C <sub>7</sub>	1.893	1.891
Si–C <sub>8</sub>	1.900	1.900
Si–C <sub>10</sub>	1.899	1.898
C <sub>1</sub> –C <sub>2</sub>	1.404	1.434
C <sub>2</sub> –C <sub>3</sub>	1.399	1.430
C <sub>3</sub> –C <sub>4</sub>	1.398	1.422
C <sub>4</sub> –C <sub>5</sub>	1.401	1.420
C <sub>5</sub> –C <sub>6</sub>	1.396	1.431
C <sub>6</sub> –C <sub>1</sub>	1.404	1.421
C <sub>8</sub> –C <sub>9</sub>	1.538	1.538
C <sub>10</sub> –C <sub>11</sub>	1.539	1.539

ular dihydrogen bond. Upon electronic excitation to the S<sub>1</sub> state of the dihydrogen-bonded phenol–DEMS complex, the Si–H<sub>2</sub> bond lengthens by only 0.004 Å and the Si–C<sub>7,8,10</sub> bonds hardly change. Thus the Si–H stretching vibrational absorption in the S<sub>1</sub> state is still observed in Fig. 3, but a stretching vibrational mode in the region of a carbonyl group is not. The distance H<sub>1</sub>···H<sub>2</sub> in the dihydrogen-bonded phenol–DEMS complex shortens significantly by 0.129 Å in the S<sub>1</sub> state compared with the ground state. This finding is consistent with the results of the calculated IR spectra which implied that the dihydrogen bond strengthens upon excitation to the S<sub>1</sub> state. The O–H<sub>1</sub>···H<sub>2</sub> and H<sub>1</sub>···H<sub>2</sub>–Si bond angles increase and decrease by only +0.7° and –6.3°, respectively, between the ground and the S<sub>1</sub> state. However, the dihedral angle O–H<sub>1</sub>···H<sub>2</sub>–Si increases significantly upon excitation from the ground state (46.65°) to the S<sub>1</sub> state (76.32°). The dihedral angle of the S<sub>1</sub> state results in the phenol and DEMS molecules becoming almost perpendicular to each other. The variations of the bond angles and dihedral angles between the S<sub>0</sub> state and the S<sub>1</sub> state are attributed to conformational changes in the complex upon electronic excitation. However, the conformation of the dihydrogen-bonded phenol–DEMS complex does not change as much as that of the dihydrogen-bonded phenol–BTMA complex upon excitation from the S<sub>0</sub> state to the S<sub>1</sub> state. Thus, it is postulated that the dihydrogen bond O–H···H–B is stronger than O–H···H–Si in the S<sub>1</sub> state of the respective complexes. Furthermore, all the C–C bonds of the phenol benzene ring lengthen upon electronic excitation, while the Si–C<sub>7,8,10</sub> and C–C bonds of the DEMS moiety are almost unchanged. The conformational changes of the benzene ring can be attributed to the LE state on the phenol moiety.

#### 4. Conclusion

In this study, the electronic excited state properties of a dihydrogen-bonded phenol–DEMS complex were investigated by TDDFT. An intermolecular dihydrogen bond O–H···H–Si existed between the phenol and DEMS molecules in the optimized geometry of the phenol–DEMS complex. According to the calculated electronic transition energies, the electronic transitions of the dihydrogen-bonded phenol–DEMS complex are blue-shifted compared with that of isolated phenol as a result of the formation of the intermolecular dihydrogen bond. Analysis of the frontier MOs revealed that the S<sub>1</sub> state of the dihydrogen-bonded phenol–DEMS complex is a LE state centered on the phenol moiety.

The dihydrogen-bonded phenol–BTMA complex also exhibits a LE S<sub>1</sub> state. The IR spectra of the S<sub>0</sub> and S<sub>1</sub> states of the dihydrogen-bonded phenol–DEMS complex were also calculated. Vibrational absorption bands associated with the O–H and Si–H bonds of the dihydrogen-bonded phenol–DEMS complex were slightly red shifted upon excitation to the S<sub>1</sub> state. Moreover, the O–H<sub>1</sub> and Si–H<sub>2</sub> bonds were both longer in the S<sub>1</sub> state relative to those in the ground state. However, the C<sub>1</sub>–O bond shortens slightly and the Si–C<sub>7,8,10</sub> bonds remain almost unchanged upon electronic excitation. Thus, the excited state IR spectra of the dihydrogen-bonded phenol–DEMS complex exhibit strong, sharp peaks for the O–H and Si–H stretching vibrational modes, and do not show C=O character. The C–C bonds of the benzene ring increase in length because of the LE state centered on the phenol moiety. The H<sub>1</sub>···H<sub>2</sub> distance of the dihydrogen bond is much shorter in the S<sub>1</sub> state than in the ground state. It can be concluded that the intermolecular dihydrogen bond O–H···H–Si of the dihydrogen-bonded phenol–DEMS complex is stronger in the S<sub>1</sub> excited state. In addition, the dihydrogen bond is stronger in a phenol–borane–amine complex than in the corresponding phenol–silane complex both in the ground and electronically excited state. Understanding the excited state properties of the dihydrogen bond continues as an area of intense experimental and theoretical research.

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